$C_n N^+$ - He Collisions – Fundamental aspects and application to astrochemistry.



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In this work, we try to understand the process of excitation and fragmentation of $C_n N^+$ molecules in high velocity collisions with atoms. Experiments are conducted at the Tandem accelerator, Orsay. The motivation being the presence of such species in astrophysical environments, for instance the atmosphere of Titan. The collision is divided into two phases, excitation of the molecule in the first phase $\{10^{-16} \text{ s}\}$ then relaxation by fragmentation. The excitation phase will be modelled theoretically in collaboration with theoreticians in Madrid and Bordeaux. The data of fragmentation are relevant in the astrochemical domain and will be added to the recent international kinetic database KIDA where researchers all over the world will use the data.

Experimental set-up and methods





Fig 2. 1d – Amplitude vs Number of counts.

Fig 3. Two dimensional representation.

Cross section for anionic fragment production

There is a competition between single and double collisions. The single collision condition is

$$P = \sigma B_{jet}$$

But anionic production probabilities are found to be quadratic with B_{jet} . This leads to the conclusion that there is indeed not one but two collisions for their formations. In case of anionic production, probability of formation is taken as

(4)

$$P = \sigma_1 B_{jet} + \frac{\sigma_2 \sigma_3 B_{jet}^2}{2}$$
(5)

which is a polynomial equation of second order with respect to B_{iet}. The experimental data agrees with this prediction given in fig. 8. On dividing (4) with a probability of species following a single collision condition P_{norm} , we get the following

2.2e-5

2.0e-5

1.8e-5

$$\frac{P}{P_{norm}} = \frac{\sigma_1}{\sigma_{norm}} + \frac{\sigma_2 \sigma_3}{2\sigma_{norm}} B_{jet}$$
(6)
which is a 1st order polynomial equation with intercept in Fig. 9 gives
and the slope gives, $\frac{\sigma_2 \sigma_3}{2\sigma_{norm}}$.

 $\sigma_{\scriptscriptstyle{
m norm}}$

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C+N

Fig 4. Silicon detectors

Fig 5. 1d – Energy vs Number of counts.

The analysis of the data from the silicon detectors Fig. 4 is done using a software DP2 [1] which is a 2D visualization and analysis software for multi-dimensional data. We design grids around patches Fig. 3 of interest. The one dimensional plots featuring amplitude and integral of the current signals are easily treated in 2d DP2 representation.

Absolute cross sections

relative positions [2].







which is a

2.5e-7

2.0e-7

Fig. 8 Evolution of P(C-/N++) with target thickness B_{iet}

CN+ analysis results

Process (σ_x)	Cross section CN ⁺ (relative error in %) (cm ²)
$\sigma_{\rm DC}$ $\rm CN^+ \rightarrow \rm CN^-$	3.9 10-20 (21%)
$\sigma_{\rm SC} \ {\rm CN}^+ \rightarrow {\rm CN}$	4.15 10 ⁻¹⁷ (17%)
$\sigma_{\rm ED} \ {\rm CN^+} \rightarrow {\rm CN^+}$	9.44 10-17 (18%)
σ_{SI} CN ⁺ \rightarrow CN ⁺⁺	2.29 10 ⁻¹⁶ (19%)
$\sigma_{\rm DI}$ CN ⁺ \rightarrow CN ⁺⁺⁺	5.44 10-17 (18%)
σ_{TI} CN ⁺ \rightarrow CN ⁺⁺⁺⁺	8.28 10-18 (18%)

Table 1. Measured cross sections of various processes(preliminary results).



Fig. 9 Evolution of P/Pnorm(C-/N++) with target thickness B_{iet}

Process of diatomic(AB+)1.	Relaxation channels	BR 2.2 a.u(relative error in %)	BR 3.6 a.u(CN+)[3]
BR(DC)	C-/N	0.02(201%)	
$CN^+ \rightarrow \{CN^-\}$	CN-	0.98(29%)	
BR(SC)	C-/N+	0.001(27%)	
$CN^+ \rightarrow \{CN\}$	C + N	0.55(1%)	0.5(1 %)
	CN	0.45(1%)	0.5(1 %)
BR(ED)	C-/N++	0.00001(29%)	
$CN^+ \rightarrow \{CN^+\}^*$	C+/N	0.55(5%)	0.53(6%)
	C/N^+	0.45(5%)	0.47(6 %)
BR(SI)	C+/N+	0.92(1%)	0.94(4 %)
$CN^+ \rightarrow CN^{++}$	C++/N	0.052(1%)	0.04(50 %)
	C/N++	0.024(10%)	0.02(100 %)
BR(DI)	C++/N+	0.58(8%)	0.61(5 %)
$CN^+ \rightarrow CN^{+++}$	N++/C+	0.40(10%)	0.39(5 %)
	N+++/C	1.3 10 ⁻³ (15%)	1 10-3 (100%)

The cross section σ of an inelastic process is calculated by the following. For single collision the equations used [2] are

$$\sigma = \frac{prob(y_0, z_0)}{B_{jet}(y_0, z_0)} \qquad (1) \qquad \frac{prob(y_0, z_0)}{\int prob(y_0, z_0) dy_0} = \frac{B_{jet}(y_0, z_0)}{\int B_{jet}(y_0, z_0) dy_0} \qquad (2) \qquad \int B_{jet}(y_0, z_0) dy_0 = \frac{dN}{dt} \frac{1}{\langle v_z \rangle} \qquad (3)$$

The jet is scanned laterally in Fig. 6 along y, the probability of collision is measured at each point and prob(y, z0) at fixed z0 is experimentally deduced. It can be shown that the denominator in Eq. (2) is related to the total flow rate of jet $\frac{dN}{dN}$ which is measured and mean velocity of the molecules in the jet $\langle v_z \rangle$ along z. dt

Future.

	C+++/N	4.0 10 ⁻³ (15%)	4 10-3(25 %)
BR(TI)	C+++/N+	0.1(3%)	0.15(6%)
$CN^+ \rightarrow CN^{++++}$	C++/N++	0.86(1%)	0.80(5 %)
	N+++/C+	0.04(3%)	0.05(20 %)

Table 2. Measured branching ratios and comparison with velocities 2.2 a.u and 3.6 a.u \rightarrow Completing the analysis for other C_nN⁺. (preliminary results) \rightarrow Another experiment in autumn 2016 with CCD (Charge Coupled Device) camera. → Modelling of the collisions CTMC (Classical Trajectory Monte Carlo) simulation with Clara Illescas in Madrid.

References

[1] L. Tassan-Got, personal communication, October 2015. [2] K. Wohrer *et al.*, Rev. Sci. instrum., **71**, 5 (2000). [3] A. Jallat (2015) PhD thesis.