

Compte rendu examen Thermo Statistique  
Decembre 2008

I) Rendement d'une lampe

1) corps noir : cavité fermée contenant un gaz de photons en équilibre thermique avec un thermostat

un corps noir absorbe intégralement le rayonnement qu'il reçoit et ne réfléchit rien.

$$2) u(\lambda, T) = \frac{8\pi hc}{\lambda^5} \frac{1}{e^{hc/\lambda k_B T} - 1}$$

$$x = \frac{hc}{\lambda k_B T}$$

$$\lambda = \frac{hc}{k_B T x}$$

$$d\lambda = \frac{hc}{k_B T} \frac{dx}{x^2}$$

$$u(T) = \frac{8\pi (hc)^5}{(hc)^5} \int_0^\infty \frac{x^3 dx}{e^x - 1} \times (k_B T)^4$$

$$u_{\text{totale}}(T) = \frac{8\pi}{h^3 c^3} (k_B T)^4 \int_0^\infty \frac{x^3}{e^x - 1} dx$$

$$E(T) = V u_{\text{totale}} = \frac{8\pi V (k_B T)^4}{h^3 c^3} \frac{\pi^4}{15}$$

$$r = \frac{E_V(T)}{E(T)} = \frac{V \int_{\lambda_1}^{\lambda_2} u(\lambda, T) d\lambda}{E(T)}$$

$$r = \frac{\int_{x_1}^{x_2} \frac{x^3}{e^x - 1} dx}{\int_0^\infty \frac{x^3}{e^x - 1} dx}$$

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$$x_1 = \frac{hc}{\lambda_1 k_B T} \quad x_2 = \frac{hc}{\lambda_2 k_B T}$$

$$r = \frac{15}{\pi^4} \int_{x_1}^{x_2} \frac{x^3}{e^x - 1} dx \quad T = 2500 \text{ K}$$

$$x_1 = 7,2$$

$$x_2 = 14,5$$

$$r = \frac{15}{\pi^4} \times \left[ -e^{-x_2} (x_2^3 + 3x_2^2 + 6x_2 - 6) + e^{-x_1} (x_1^3 + 3x_1^2 + 6x_1 - 6) \right]$$

$$3) r \uparrow \text{ sur } T \uparrow \quad (-0,019 + 0,431) \times 15/\pi^4 = 0,066$$

## II Vibrations d'un solide - Modèle d'Einstein

$$1) \quad Z = \sum_{n_x, n_y, n_z=0}^{\infty} \exp(-\beta \epsilon_{n_x, n_y, n_z}) = \sum_{n_x, n_y, n_z=0}^{\infty} \exp(-\beta \left[ (n_x + \frac{1}{2}) + (n_y + \frac{1}{2}) + (n_z + \frac{1}{2}) \right] \hbar \omega}$$

$$Z = \sum_{n_x=0}^{\infty} \exp(-\beta \epsilon_{n_x}) \sum_{n_y=0}^{\infty} \exp(-\beta \epsilon_{n_y}) \sum_{n_z=0}^{\infty} \exp(-\beta \epsilon_{n_z})$$

$$Z = \left( \sum_{n_x=0}^{\infty} \exp(-\beta (n_x + \frac{1}{2}) \hbar \omega) \right)^3$$

$$Z = \left[ \exp(-\beta \frac{\hbar \omega}{2}) \sum_{n_x=0}^{\infty} \exp(-\beta n_x \hbar \omega) \right]^3$$

$$\sum_{n_x=0}^{\infty} \exp(-\beta n_x \hbar \omega) = 1 + \exp(-\beta \hbar \omega) + \exp(-2\beta \hbar \omega) + \dots$$

→ somme d'une série géométrique de raison  $\exp(-\beta \hbar \omega)$

$$Z = \left( e^{-\beta \frac{\hbar \omega}{2}} \right)^3 \left( \frac{1}{1 - e^{-\beta \hbar \omega}} \right)^3$$

$$Z = Z^N = \left( e^{-\beta \frac{\hbar \omega}{2}} \right)^{3N} \left( \frac{1}{1 - e^{-\beta \hbar \omega}} \right)^{3N}$$

$$Z = e^{-\beta \frac{3}{2} N \hbar \omega} \frac{1}{(1 - e^{-\beta \hbar \omega})^{3N}}$$

$$\langle E \rangle = - \frac{\partial}{\partial \beta} \ln Z = \frac{3}{2} \hbar \omega + \frac{3 \hbar \omega e^{-\beta \hbar \omega}}{1 - e^{-\beta \hbar \omega}}$$

$$\langle E \rangle = \frac{3}{2} \hbar \omega + \frac{3 \hbar \omega}{e^{\beta \hbar \omega} - 1}$$

$$\langle E \rangle = N \langle E \rangle$$

$$F = -k_B T \ln Z = N \left[ \frac{3}{2} \hbar \omega + 3 k_B T \ln(1 - e^{-\beta \hbar \omega}) \right]$$

$$C_V = \left( \frac{\partial \langle E \rangle}{\partial T} \right)_V = \left( \frac{\partial \langle E \rangle}{\partial \beta} \right)_V \frac{\partial \beta}{\partial T} = - \frac{1}{k_B T^2} \left( \frac{\partial \langle E \rangle}{\partial \beta} \right)_V$$

$$C_V = \frac{3N \hbar \omega}{k_B T^2} \frac{\hbar \omega e^{+\beta \hbar \omega}}{(e^{+\beta \hbar \omega} - 1)^2} = 3N k_B \left( \frac{\hbar \omega}{k_B T} \right)^2 \frac{e^{\beta \hbar \omega}}{(e^{\beta \hbar \omega} - 1)^2}$$

$$\theta_E = \frac{h\nu}{k_B}$$

$$C_V = 3Nk_B \left(\frac{\theta_E}{T}\right)^2 \frac{e^{+\theta_E/T}}{(e^{\theta_E/T} - 1)^2}$$

$T > \theta_E$  Hautes températures  $\frac{\theta_E}{T} < 1$

$$e^{\theta_E/T} \sim 1 + \frac{\theta_E}{T}$$

$$e^{\theta_E/T} - 1 \sim \frac{\theta_E}{T}$$

$$C_V \sim 3Nk_B \left(\frac{\theta_E}{T}\right)^2 \frac{(1 + \theta_E/T)}{(\theta_E/T)^2}$$

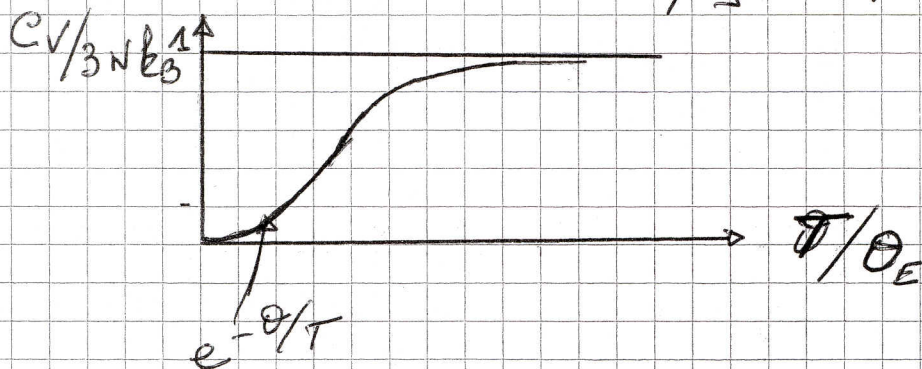
$$C_V \sim 3Nk_B \cancel{\left(\frac{\theta_E}{T}\right)^2} = \cancel{3Nk_B} \sim 3Nk_B$$

$T < \theta_E$  basses températures  $\frac{\theta_E}{T} > 1$

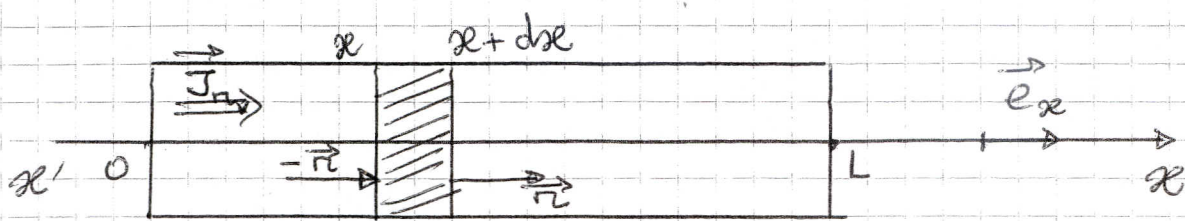
$$e^{\theta_E/T} - 1 \sim e^{\theta_E/T}$$

$$C_V \sim 3Nk_B \left(\frac{\theta_E}{T}\right)^2 e^{-\theta_E/T}$$

tend vers 0 avec T de façon exponentielle



### III Diffusion axiale de molécules



1) Loi de Fick

$$\vec{J}_{n_v} = -D \left( \frac{dn_v(x)}{dx} \right) \vec{e}_x$$

$\vec{J}_{n_v}$  : vecteur courant volumique de molécules en  $m^{-2}s^{-1}$   
 $D$  : coefficient de diffusion en  $m^2s^{-1}$   
 $\frac{dn_v(x)}{dx}$  : gradient de densité en  $m^{-4}$

2)  $n_v(x=0) = n_1$        $n_v(x=L) = n_2$

$$dN^V = dt \iint_S \vec{J}_{n_v} \cdot (-\vec{n}) ds \quad dN^P = 0$$

$$(dn_v) S dx = \left[ S J_{n_v}(x) - S J_{n_v}(x+dx) \right] dt$$

$$= - \left( \frac{\partial J_{n_v}}{\partial x} \right) dx dt$$

$$\frac{dn_v}{dt} = - \frac{\partial J_{n_v}}{\partial x} = - \text{div} \vec{J}_{n_v}$$

3)  $\frac{dn_v}{dt} = - \frac{\partial}{\partial x} \left( -D \frac{dn_v}{dx} \right) = -D \frac{d^2 n_v}{dx^2} = 0$

$$\frac{d^2 n_v}{dx^2} = 0$$

4)  $\frac{dn_v}{dx} = cte = - \frac{J_{n_v}}{D}$        $n_v(x) = - \frac{J_{n_v}}{D} x + cte$

$$n_v(0) = cte = n_1$$

$$n_v(L) = - \frac{J_{n_v}}{D} L + cte$$

$$n_v(L) = - \frac{J_{n_v}}{D} L + n_v(0)$$

$$n_2 - n_1 = - \frac{J_{n_v}}{D} L$$

$$\frac{n_2 - n_1}{L} = - \frac{J_{n_v}}{D}$$

$$n_v(x) = \frac{(n_2 - n_1)}{L} x + n_1$$

$u(\lambda, T)$

