

Mouvement sans l'action de la pesanteur:

$$1. \Sigma \vec{F} = m \vec{a}_{AIR} \rightarrow \begin{cases} m\ddot{x} = 0 \\ m\ddot{y} = 0 \\ m\ddot{z} = -mg \end{cases} \rightarrow \begin{cases} \dot{x} = C_0 = 0 \\ \dot{y} = C_0 = v_0 \cos \alpha \\ \dot{z} = -gt + C_0 = -gt + v_0 \sin \alpha \end{cases}$$

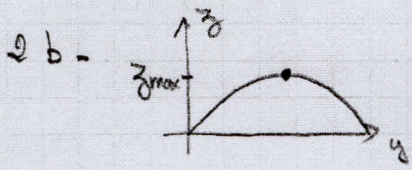
$$\hookrightarrow \begin{cases} x = C_0 = 0 \\ y = v_0 \cos \alpha t + C_0 = v_0 \cos \alpha t \\ z = -\frac{1}{2}gt^2 + v_0 \sin \alpha t + \frac{C_0}{g} \end{cases}$$

$x = C_0 \Rightarrow$ mouvement plan

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2 a. eq. cartésienne: $z(y) \Rightarrow z = -\frac{g}{2v_0^2 \cos^2 \alpha} y^2 + \tan \alpha y \rightarrow$ parabole

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2 b. z_{max} tel que $\frac{dz}{dy} \Big|_{z_{max}} = 0$

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ou simplement tel que \dot{z} en $z_{max} = 0$

Après calcul: $z_m = \frac{1}{2g} \sin^2 \alpha \cdot v_0^2$

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2 c. $y = y_c$ pour $z = 0 \Rightarrow y \left(-\frac{g}{2v_0^2 \cos^2 \alpha} y + \tan \alpha \right) = 0$

$\hookrightarrow y_c = -\tan \alpha \times -\frac{2v_0^2 \cos^2 \alpha}{g} = \frac{2v_0^2}{g} \sin \alpha \cos \alpha$

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α_{max} tel que $\frac{dy_c}{d\alpha} \Big|_{\alpha_{max}} = 0$

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$\hookrightarrow \frac{dy_c}{d\alpha} = \frac{2v_0^2}{g} (\cos^2 \alpha - \sin^2 \alpha) = 0$ pour $\cos \alpha = \sin \alpha$
 $\alpha = \frac{\pi}{4}$

3. $\Sigma \vec{F} = m \vec{a}_{AIR} \rightarrow \begin{cases} m\ddot{x} = -kx & (a) \\ m\ddot{y} = -ky & (b) \\ m\ddot{z} = -mg - kz & (c) \end{cases}$

$\hookrightarrow (a) \ddot{x} + \frac{k}{m} x = 0 \rightarrow \dot{x} + \frac{k}{m} x = 0 \rightarrow x = A e^{-\frac{k}{m}t}$
C.I. $\Rightarrow A = 0$

(b) $\ddot{y} + \frac{k}{m} y = 0 \rightarrow \dot{y} + \frac{k}{m} y = 0 \rightarrow y = B e^{-\frac{k}{m}t}$
C.I. $\Rightarrow B = v_0 \cos \alpha$

(c) $\ddot{z} + \frac{k}{m} z = -g \rightarrow z = C e^{-\frac{k}{m}t} - \frac{gm}{k}$
C.I. $\Rightarrow C = v_0 \sin \alpha + \frac{gm}{k}$

$\Rightarrow \begin{cases} \dot{x}(t) = 0 \\ \dot{y}(t) = v_0 \cos \alpha e^{-\frac{k}{m}t} \\ \dot{z}(t) = \left(v_0 \sin \alpha + \frac{gm}{k} \right) e^{-\frac{k}{m}t} - \frac{gm}{k} \end{cases}$

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\hookrightarrow qd $t \rightarrow +\infty, \vec{v}_{AIR} \rightarrow \begin{pmatrix} 0 \\ 0 \\ -\frac{gm}{k} \end{pmatrix}$

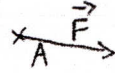
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Mécanique du solide

Questions de cours:

1. $R^* = (C, \vec{e}_x, \vec{e}_y, \vec{e}_z)$ où C est le centre de masse du solide considéré

2. $\vec{M}_O = \vec{OA} \wedge \vec{F}$



O x

3. $\sum \vec{F}_{ext} = m \vec{a}_{C/R}$

$\sum \vec{M}_O(\vec{F}_{ext}) = \frac{d\vec{L}_O}{dt} \Big|_R$

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Éléments cinétique d'une boule de billard:

1a. $C_x, C_y, C_z \rightarrow [I]_C = \begin{pmatrix} I_{Cx} & 0 & 0 \\ 0 & I_{Cy} & 0 \\ 0 & 0 & I_{Cz} \end{pmatrix}$

b. $I_C = \int_{vol} (x^{*2} + y^{*2} + z^{*2}) \rho dV = \int_{vol} r^2 \rho dV = \frac{M}{\frac{4}{3}\pi a^3} \iiint r^2 dxdydz = \frac{3}{5} Ma^2$

c. $I_{Cx} = \int_V (y^{*2} + z^{*2}) \rho dV$; $I_{Cy} = -$; $I_{Cz} = -$

$\left\{ \begin{array}{l} I_{Cx} + I_{Cy} + I_{Cz} = 2I_C \\ I_{Cx} = I_{Cy} = I_{Cz} \end{array} \right.$

$\rightarrow I_{Cx} = I_{Cy} = I_{Cz} = \frac{2}{3} I_C$

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2a. $\vec{L}_{C/R} = [I]_C \cdot \vec{\omega} = \omega I_{Cx} \vec{e}_x$

$E_{k/R^*} = \frac{1}{2} \vec{L}^* \cdot \vec{\omega} = \frac{1}{2} \vec{L}_{C/R} \cdot \vec{\omega} = \frac{1}{2} \omega^2 I_{Cx}$

$\omega = \frac{v}{r} = \frac{v}{\frac{2}{3} r}$

b. $CRSG \Rightarrow \vec{v}_g = \vec{0} \Rightarrow \vec{v}_{I_1/R} = \vec{v}_{I_2/R} \Rightarrow \vec{v}_{C/R} + (\vec{\omega} \wedge \vec{r}_{C/I_1}) = \vec{0}$

$\rightarrow v_0 + \omega a = 0$

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