

# EXPERIMENTAL STUDIES OF DISSOCIATIVE RECOMBINATION OF $H_3^+$ , $KrH^+$ AND $XeH^+$ .

A.P. Le Padellec, S. Laubé, C. Rebrion,  
O. Sidko, J.C. Gomet and B.R. Rowe.

*Département de Physique Atomique et Moléculaire, URA 1203 du CNRS,  
Université de Rennes I, Campus de Beaulieu,  
35042 Rennes Cedex, France.*

The exact value of the rate coefficient  $\alpha$  for dissociative recombination of  $H_3^+$  ground state ions has been and is still a subject of controversy<sup>(1),(2),(3)</sup>. Following our previous work in helium afterglow, we have measured  $\alpha(H_3^+)$ ,  $\alpha(KrH^+)$  and  $\alpha(XeH^+)$  in an argon-helium buffer using a FALP-MS apparatus.

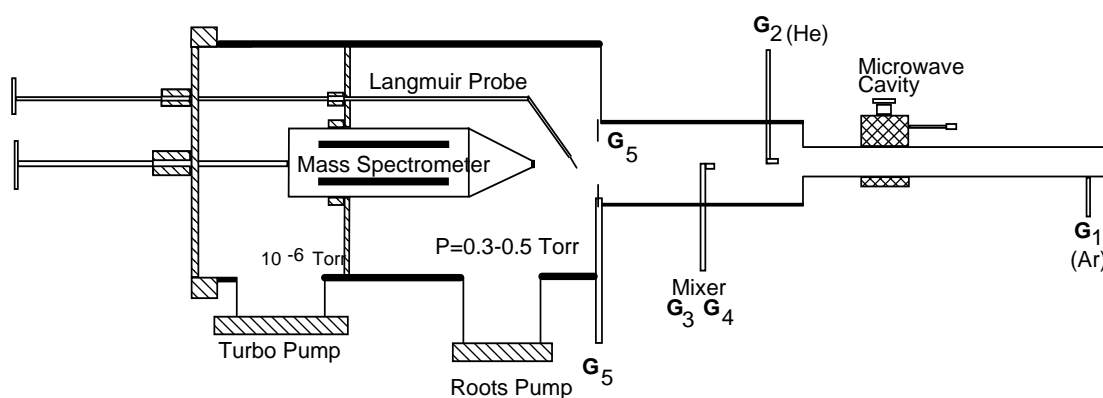
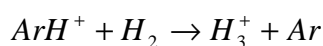
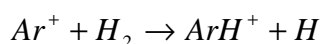


Fig.1:

Figure 1: MS-FALP experimental set-up

Argon flows through the discharge (Gate  $G_1$ :  $Q_{Ar}=30 \text{ sl min}^{-1} \text{ atm}$ ) and the electrons are thermalized downstream by large helium injection ( $G_2$ :  $Q_{He}=5 \text{ sl min}^{-1} \text{ atm}$ ). It is therefore possible to obtain at low pressure ( $P \sim 0.5 \text{ Torr}$ ) a plasma where  $Ar^+$  is the dominant ion. Addition of hydrogen alone results in  $H_3^+$  production through the following reactions:



The latter reaction is known to form  $H_3^+$  up to  $v=2$ . In this case,  $\alpha(H_3^+)$  has been determined to be  $1.0 \cdot 10^{-7} \text{ cm}^3 \text{ s}^{-1}$ .

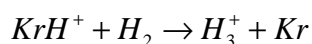
When Kr or Xe (Gate  $G_3$ ) is injected in excess of  $H_2$ , (Gate  $G_4$ ) the plasma is dominated by  $KrH^+$  or  $XeH^+$ , these two gazes having proton affinities larger than  $H_2$ . The measured values of the rate coefficients for these ions are respectively:

$$\alpha(KrH^+) < 1.0 \cdot 10^{-8} \text{ cm}^3 \text{ s}^{-1}$$

$$\alpha(XeH^+) = 1.1 \cdot 10^{-7} \text{ cm}^3 \text{ s}^{-1}$$

in good agreement with the previous work of Geoghegan<sup>(4)</sup> and coworkers only for  $KrH^+$ .

By adding a large quantity of  $H_2$  further downstream (Gate  $G_5$ ) in a flowing afterglow plasma dominated by  $KrH^+$ , it is possible to obtain  $H_3^+$  as a dominant ion through the reaction:



with  $[H_2] \gg [Kr]$  due to the very close proton affinity of Kr and  $H_2$ . Since  $KrH^+$  does not recombine, this can be done with a fairly high density. It is therefore possible to measure  $\alpha(H_3^+)$  for ions that are almost certainly in their ground state: a computer simulation supports this conclusion.

We have in our disposal two different methods to calculate  $\alpha$  (see text below). Depending of the method used, we got :

$$\alpha(H_3^+) = 0.9 \cdot 10^{-7} \text{ cm}^3 \text{ s}^{-1} \quad (\text{Method 1: figure 2})$$

$$\alpha(H_3^+) = 0.8 \cdot 10^{-7} \text{ cm}^3 \text{ s}^{-1} \quad (\text{Method 2: figure 3})$$

Considering the uncertainties in the various experiments, this result is in good agreement with the measurements of Larsson and coworkers<sup>(5)</sup>. However it can be taken as showing that the  $\alpha$  value for ground state is slightly lower than for excited states.

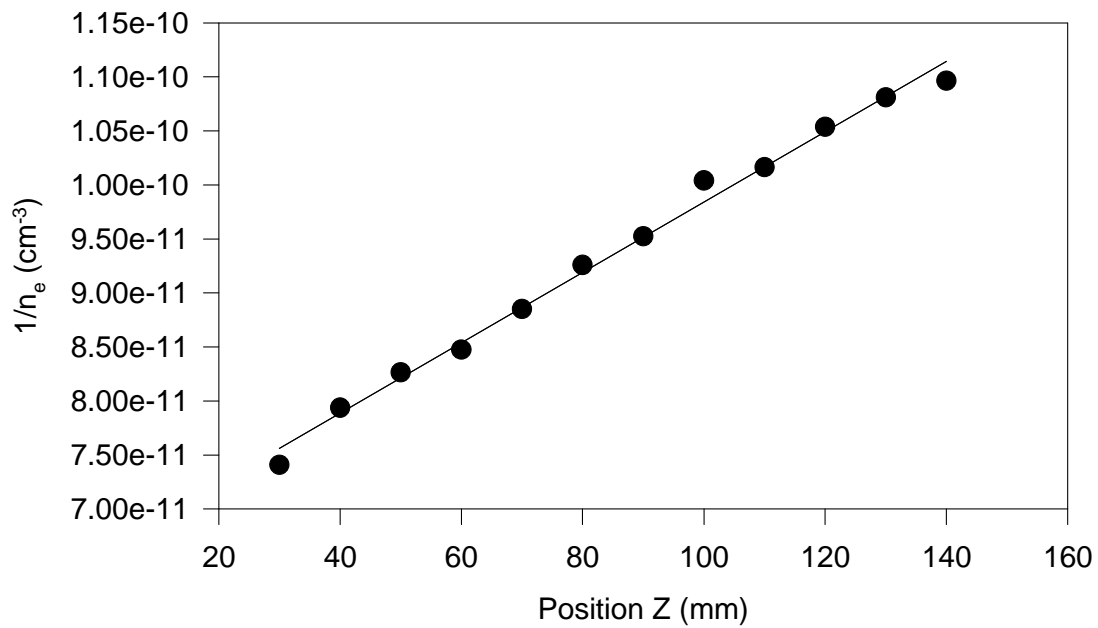


Figure 2: First Method  $\alpha(H_3^+) = 0.9 \times 10^{-7} \text{ cm}^3 \text{ s}^{-1}$

We dispose of two different methods to deduce the rate coefficient.

1<sup>st</sup> Method: The slope of the plot  $\frac{1}{n_e}$  versus  $Z$  yields the ratio  $\alpha/v$ . This method does not take into account destruction processes of  $H_3^+$  other than dissociative recombination.

2<sup>nd</sup> Method: For a fixed position  $Z$  relative to a position of reference  $Z_0$ , the slope of the plot  $\ln \frac{[H_3^+]_Z}{[H_3^+]_{Z_0}}$  versus  $\frac{1}{v} \int_{Z_0}^Z n_e dz$  yields the rate coefficient. This method takes into account the destruction of  $H_3^+$  by ion-molecule reactions. The only assumption is that the molecule density is constant over the whole range of position  $Z$ .

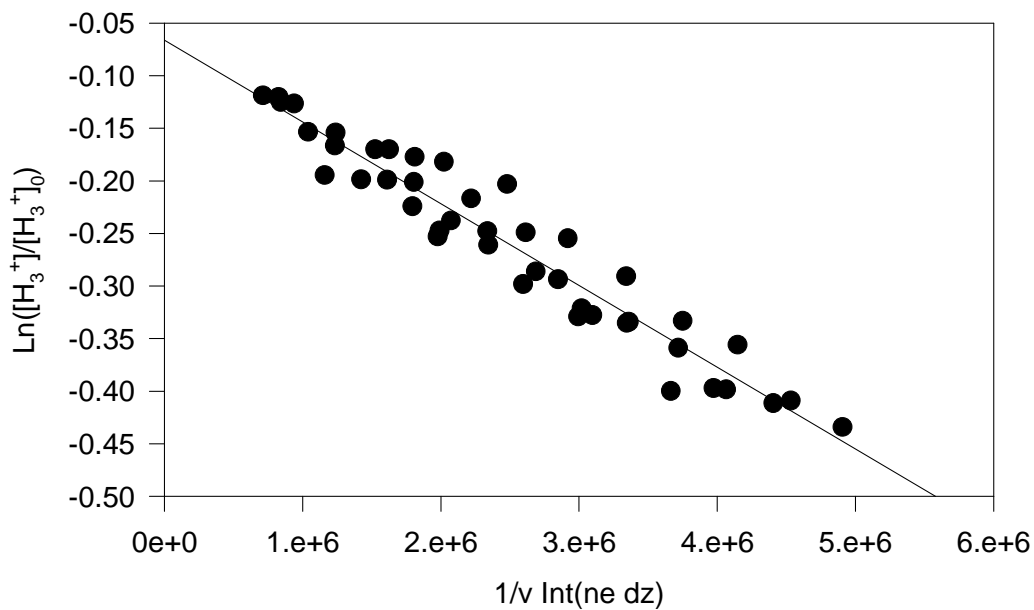


Figure 3: Second Method  $\alpha(\text{H}_3^+) = 0.8 \times 10^{-7} \text{ cm}^3 \text{ s}^{-1}$

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